APPLICATIONS OF A NODAL-INTEGRATION-BASED FINITE ELEMENT METHOD TO NON-LINEAR PROBLEMS

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Abstract. In this paper, we firstly introduce a nodal-integration-based finite element method. The method allows the use of first-order tetrahedral elements without suffering from the volumetric locking problem. The most important advantage of tetrahedral meshes is that they can be automatically generated for complex geometries using existing reliable meshing tools. The method is then applied to 3 types of applications. The first application is a large displacement, large strains elastic-plastic simulation on a notched specimen. The second application is an elastic-plastic bending problem. And the last example concerns the numerical simulation of the thermo-mechanical problem. In all the cases, the solution given by the nodal-integration-based FEM is compared to more classical FEM results.

1 INTRODUCTION

Finite element simulations of the behavior of structures made up of materials obeying the von Mises plasticity criterion (the most commonly used criterion especially for metal materials) face particular difficulties. These difficulties come from the plastic incompressibility condition imposed by the behavior of the material. This condition of (nearly-)incompressibility which should be resolved at each integration point can lead to a volumetric locking phenomenon.

Finite element formulations have thus been proposed to alleviate these difficulties. These formulations are generally based either on reduced numerical integration schemes, or on mixed formulations of the problem [1, 2]. In practice, formulations based on reduced integration schemes are easily applicable only with hexahedral elements. The absence of automatic meshing tools with this type of finite element then leads engineers to tedious operations which are costly in human time. The existence of automatic tetrahedral meshing tools gives this type of finite element an indisputable economic advantage. New

first order tetrahedral finite elements have thus been developed in recent years on the basis of mixed displacement-pressure formulations of the problem [3, 4, 5]. However, the presence of different degrees of freedom (displacement and pressure) or even that of internal degrees of freedom for certain elements, such as the element P1+P1 [6, 7], can then lead to significant additional computation costs.

In the literature, choosing the nodes as integration points can be found both in meshless methods and in the finite element method. In meshless methods especially, the Stabilized Conforming Nodal Integration (SCNI) technique proposed by Chen *et al.* [8, 9, 10] performs very well in several benchmark problems. The principle of SCNI technique is to consider nodal deformations calculated on average over a volume surrounding the node. The volume integral is then transformed into a surface integral thanks to the divergence theorem (Figure 1). These volumes can be obtained for example from a tessellation of Voronoi cell. Puso *et al.* [11] has proposed a schema of improvement for solving low energy mode that appears in the nodal integration technique.

Figure 1: Cell integration for SCNI [16].

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As far as the finite element method is concerned, Bonnet *et al.* [12] has proposed an average nodal pressure tetrahedral element for explicit dynamic applications. The element avoids volumetric locking problems by computing average nodal pressures from surrounding elements. Dohrmann *et al.* [13] *et al.* presented a new tetrahedral element for small strain elasticity with an averaging procedure now applied on all the strain components. The formulation prevents the volumetric locking by producing a favorable constraint ratio for the volumetric response. Bonnet *et al.* [14] has extended this node-based formulation to explicit dynamics applications. Krysl *et al.* [15] presents assumed-strain finite elements based on nodal integration. They show that the three-node triangle, four-node tetrahedron, and eight-node and 27-node hexahedra give good performance in both compressible and almost incompressible regime. The combined use of the Finite Element Method with the SCNI technique has been developed in 2D on the basis of triangular meshes by Qual *et al.* [16] and Canales *et al.* [17]. Ce method has been extended for 3D thermo-mechanical applications by Jia *et al.* [18]. In this paper, we present several benchmark tests by comparing the results obtained with the nodal-integration-based finite element method as proposed by Jia *et al.* [18] with those coming from more classical finite element solutions. All the simulations are performed using *SYSWELDTM* software.

2 Notched tensile specimen

The objective of the first test is to compare the nodal-integration-based solution to classical FEM solutions on a 3D large strain elastoplastic problem. For symmetry reasons, only half of the specimen on an angular area is considered. Symmetry boundary conditions are applied to the faces of the model. The numerical model is presented in figure 2. Four numerical simulations based on the following formulations have been performed and compared:

1. P1 element: linear variation of the displacements inside the element,
2. P1P1 element: linear variation of the displacements and pressure inside the element
3. P1+P1 element: linear variation of the displacements and pressure inside the element plus 3 additional internal degrees of freedom associated with a bubble shape function.

Figure 2: Mesh of the notched tensile specimen (unit: mm).

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Large displacements, large strains are taken into account with an updated lagrangian formulation. The steel behavior is supposed to be elastoplastic with an isotropic hardening. The Young's modulus $E = 200,000$ MPa, the Poisson ratio $\nu = 0.3$ and the yield stress is a function of the cumulated equivalent plastic strain $\sigma_y(\epsilon_{eq}^p) = A(\epsilon_{eq}^p + \epsilon_{eq0}^p)^n$ with $A = 500$ MPa, $n = 0.1$ and $\epsilon_{eq0}^p = 0.01$. Initial yield stress $\sigma_{y0} = A(\epsilon_{eq0}^p)^n = 315.48$ MPa.

Figure 3 provides comparisons of force-axial displacement curves and CPU times. P1 element is known to be too stiff and to present volumetric locking problems. As expected it gives a result far from those from the other formulations. The nodal-integration-based FEM leads to a force-axial displacement curve very close to those obtained using P1P1 and P1+P1 elements, without any volumetric locking problem while requiring less CPU time.

Figure 4 shows the axial displacement and axial stress distributions obtained with P1+P1 element and the nodal-integration-based FEM. The displacement and stress distributions are very close.

Figure 3: Force-axial displacement curves and CPU times.

Figure 4: Stresses given by nodal-integration-based FEM and P1+P1 element.

3 Large deformation elastoplastic bending

As we know, the tetrahedral P1 element exhibits poor performances in bending dominated problems. Therefore the present section is dedicated to the investigation of bending performances of different approaches. We consider a beam on simple support at its ends and loaded in the middle. For symmetry reasons, only a quarter of the beam is considered. Figure 5 shows the model including loading and boundary conditions and the meshes considered. Mesh (a) is prepared for Q2 hexahedral elements. It is well known that Q2 elements give very accurate results for bending problems [1, 2]; the Q2 element solution will be therefore considered as a reference solution. Mesh (b) is prepared for Q1P0 hexahedral elements based on a mixed displacement (Q1)-pressure (P0) formulation. Mesh (c) is prepared for linear tetrahedral elements (P1P1 element, nodal-integration-based FEM). The dimensions of the model are $1 * 2 * 12 \text{ mm}$. An elastoplastic behavior law is considered with isotropic hardening and a constant

Figure 5: Numerical model and boundary conditions for bending simulations; meshes—(a) hexahedral mesh: 1,721 nodes/ 280 Q2 elements; (b) hexahedral mesh: 3,075 nodes/ 2,240 Q1 elements; (c) tetrahedral mesh: 3075 nodes/ 13,440 tetrahedral elements.

hardening slope. The material properties are given in Table I.

Table I: Material properties.

Yuong's Modulus	Poisson ratio	Yield stress	Hardening slope
$E = 195,122\text{MPa}$	$\nu = 0.30$	$\sigma_Y = 170\text{MPa}$	$H = 5,000\text{MPa}$

Figure 6 shows the contours of von Mises' stress and U_z vertical displacement. The von Mises' stress distributions obtained with the different formulations are rather similar. The nodal-integration-based FEM gives a vertical displacement distribution very close to the ones obtained using Q2 and Q1P0 elements while the P1P1 element exhibits a too stiff bending behavior. This comes from the fact that because integration points are the nodes of the mesh, the nodal-integration-based FEM leads to an underintegration of the problem.

Table II provides a comparison of the maximum vertical displacement and of the CPU times obtained with the different approaches. Compared with the reference solution (Q2 element), the nodal-integration-based FEM exhibits a behavior a little too flexible. But, despite the increase of the matrix bandwidth it leads [18], the nodal-integration-based FEM gives a CPU time significantly less than the ones obtained with Q1P0 and P1P1 elements for a problem gathering the same number of nodes. This is related to the elastoplastic resolutions that are performed at the elements integration points with the Q1P0 and P1P1 meshes and at the nodes (which are much less numerous than the integrations points) with the nodal-integration-based FEM.

Table II : Maximum displacement in Z direction.

element type	integration point	U_z	Δz	$\frac{\Delta z}{ U_z }$	CPU
Q2 element	5600	-3.95	reference	reference	187
Q1P0 element	17920	-3.88	-0.07	-0.0177	578
P1P1 element	13440	-3.38	-0.57	-0.144	860
nodal-integration-based FEM	3075	-4.11	0.16	0.04	256

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4 Thermo-mechanical simulation

In this section, we will present the application of thermo-mechanical simulation. The dimensions of problems and meshes are shown in figure 7. This problem is solved by performing a 2D axisymmetry simulation. A gaussian heat source is applied and its energy distribution is described by equation (1). The heat source parameters are $a = 8$, $b = 2$, and Q_0 which is given in Table II. The center of the heat source is defined at ($X_0 = 0$, $Y_0 = 5$). 316L material properties taken from *SYSWELDTM* database [19] is used in this simulation. The convective and radiative losses are taken into account.

$$Q_{(X,Y)} = Q_0(t) * SQRT(1 - \frac{(X - X_0)^2}{a^2} - \frac{(Y - Y_0)^2}{b^2}) \quad (1)$$

Table II: Q_0 in function of time.

time (second)	0	1.5	3	100
Q_0	0	20	0	0

Figure 7: Problem studied and meshes— (a) triangle mesh: 942 nodes / 1849 2D elements; (b) quadrilateral mesh: 1281 nodes / 1200 2D elements.

Figure 8 shows the residual stresses given by the nodal-integration-based FEM and classical FEM and the contours of residual stresses show no difference.

Figure 8: Residual stresses simulated by (a): nodal-integration-based FEM and (b): Q1 FEM.

5 CONCLUSIONS

Three applications of the nodal-integration-based finite element method have been presented. These applications are representative of different classical problems: large strain elastoplasticity, bending dominated problem, thermo-mechanical application. All the simulation are carried out with the same PC.

According to the comparisons, the nodal-integration-based FEM gives accurate results. The main advantage of the method lies on the use of tetrahedral meshes. There is no longer any volumetric locking problem. Meshing of complex structures is so greatly facilitated by the use of existing reliable meshing tools. The other benefits of nodal-integration-based FEM come from the fact that all the mechanical

quantities (displacements, strains, stresses, internal variables) are now defined at nodes. So the results files are much smaller than those coming from the classical FEM. In addition, post-processing is easier as it does not need an averaging procedure to extrapolate stresses and internal variables from Gauss points to the nodes. An easier transfer of mechanical quantities between two meshes is also expected. The main drawback of the method comes from a somewhat larger bandwidth of the tangent matrix. This drawback is compensated by the fact that the material constitutive equations are now solved at nodes that are much less numerous than Gauss points.

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