MOISTURE TRANSPORT IN POLYAMIDE (PA6) AND ITS EFFECT ON THE MECHANICAL BEHAVIOUR

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Abstract. Polyamide exhibits hygroscopic nature and can absorb up to 10 % of moisture relative to its weight. The absorbed moisture increases the mobility of the molecular chains and causes a reduction in the glass transition temperature. Thus, depending on the moisture distribution, a polyamide component can show different stiffness and relaxation times.

The moisture distribution also depends on the mechanical loading of the material. However, it was noticed that the diffusion process remains unaffected when the process is compared for a loaded and an unloaded material. It is postulated that the moisture redistribution is due to an external force that takes place as a result of the pressure acting on the moisture. The diffusion process is unaffected by preloading of the dry material as the pressure is applied purely on the material and not on the absorbed moisture. However, when a saturated specimen is loaded, the pressure is exerted on the moisture too which causes its redistribution.

In this work, the distribution of the absorbed moisture is simulated by a non-linear diffusion model. It is coupled with the viscoelastic behaviour of PA6. The stiffness of the viscoelastic model changes and the relaxation time reduces with increasing moisture concentration. The coupling of diffusion to mechanical loading is achieved with the recalculation of the moisture concentration caused due to the redistribution of volume. It is assumed that there is no transport of moisture, but the transport of volume and the change in volume creates a change in concentration in the specimen. This strongly coupled model has been implemented using the finite element method. The model results are compared to experiments for validation. A strongly coupled model was thus created which could reproduce the experimental results with reasonable accuracy.

1 INTRODUCTION

The use of polymer has increased continuously in the last decades. In the last 50 years, the consumption of polymer has increased from 5 million ton to 100 million ton worldwide [22]. Technical polymers such as polyamide (PA6) form a major portion of the polymer applications. Lightweight, low cost and manufacturability are some of the reasons why PA6 is used more often. Apart from that, its chemical resistance is useful in applications such as in the automobile industry where they are used in fuel supply components, air intake manifolds and engine covers [5, 23–25]. However, PA6 is used under various

environmental conditions, where the temperature and humidity can vary with time. As it is hygroscopic in nature, it can absorb from 9.5 % to 10 % of moisture by weight [1, 2, 9]. The diffusion of water in the polyamide depends on various parameters such as temperature, loading conditions and time. When the surrounding humidity varies polyamide can absorb water at high relative humidities but due to plastification, it can again give up some moisture causing an in-stationary and inhomogeneous distribution of moisture in the material [4, 8, 10, 16]. In the industry however, the local distribution of moisture is not taken into account and rather a saturated and a dry material is studied. The presence of moisture uptake behaviour is also effected by the loading as the free volume available for moisture absorption changes. In order to model the behaviour of PA6 under the presence of moisture, it is important to understand the

effect of water on the mechanical properties as well as the effect of mechanical loading on the diffusion process. The water molecules can form hydrogen bonds with the polymer chains [19], which causes the gliding of the chains against each other. This leads to a transformation from a glassy state to a rubbery state, thus lowering the glass transition temperature [13]. One of the effects of the phase transition is that the stiffness of the material decreases which has been seen in experimental results [6, 7, 11]. Many authors have modelled this behaviour with the help of dependency of material parameters on the concentration [3,12,14,15,18,21]. There is also a change in the viscoelastic relaxation behaviour with the phase transition, which has been captured by interpolating the relaxation times between the dry and saturated specimen to an effective relaxation time [14,21]. Further, the coupling of the diffusion model on the mechanical loading has been done by making the diffusion coefficient a function of the strain values [3, 15]. These models have been developed however on a theoretical basis and have not been verified with the help of experimental results, specially for PA6.

Experiment conducted on PA6 to study the relationship of diffusion coefficient and the strain values however gave unexpected results. It was noticed that the diffusion speed remains same for a loaded and an unloaded specimen. Hence a coupling of the diffusion coefficient to the strain is not a correct modelling approach. The model presented by the authors [21] with one sided coupling, where the moisture concentration affects the mechanical properties was to be validated with 3-point-bending tests. The results from the experiment however showed deviation from the simulation, which indicates that a moisture redistribution does take place due to the loading. Hence a coupling between the moisture content and the mechanical loading is indeed important. In the current work, the one sided coupled model has been extended and a strong coupling has been implemented, without effecting the diffusion process directly. A phenomenological approach was taken and instead of moisture redistribution the volume redistribution was considered to calculate the concentrations after loading. In the following sections a short overview of the mechanical and the diffusion model is given, and then the coupling methodology is explained. Finally the results of the fully coupled model is compared with the experimental results.

2 MODELLING

A short introduction of the modelling approach for the material behaviour and moisture uptake kinetics is given here. For a more detailed explanation the readers can refer to Sharma et al. [21].

2.1 Mechanical model

Polymers exhibit viscoelastic behaviour and PA6 is no exception. It is modelled with the help of a rheological model with four Maxwell elements (Figure 1). The elements are made up of elastic spring



Figure 1: Rheological model for viscoelastic behaviour with four Maxwell arms.

elements with Lamé parameters λ_j, μ_j for the j_{th} element. The Maxwell elements are connected to the basic elasticity spring element with Lamé parameters λ_o, μ_o . The relaxation time τ_j for each element is fixed and arranged in power of tens. The shifting of the relaxation spectrum is achieved by modelling the dry material with Maxwell elements of longer relaxation times and the saturated material with the Maxwell elements of shorter relaxation times (Table 1). The rheological model is constrained to small deformations and hence the stress σ is calculated with the help of Hooke's law given by,

$$\boldsymbol{\sigma} = 2\mu_o \boldsymbol{\varepsilon} + \lambda_o \operatorname{tr} \boldsymbol{\varepsilon} \mathbf{I} + \sum_{j=1}^4 (2\mu_j \boldsymbol{\varepsilon}_{ej} + \lambda_j \operatorname{tr} \boldsymbol{\varepsilon}_{ej} \mathbf{I}).$$
(1)

Here **I** is the identity tensor, the operator 'tr' gives the trace of the tensor and $\boldsymbol{\varepsilon}$ is the total strain which is split additively,

$$\boldsymbol{\varepsilon} = \boldsymbol{\varepsilon}_{ej} + \boldsymbol{\varepsilon}_{ij}. \tag{2}$$

 $\mathbf{\epsilon}_{ej}$ is the elastic part of the strain and $\mathbf{\epsilon}_{ij}$ is the inelastic part in the j_{th} Maxwell element. The evolution of the inelastic strains are calculated with the help of the relaxation times of the Maxwell elements with the relation,

$$\dot{\mathbf{\varepsilon}}_{ij} = \frac{\mathbf{\varepsilon} - \mathbf{\varepsilon}_{ij}}{\tau_j/2}.$$
(3)

2.2 Moisture uptake model

The moisture uptake in gravimetric experiments showed that the mass of the moisture absorbed increases linearly with the square root of time. This motivates the use of Fick's model which at higher concentration shows a deviation from the experimental results. Hence the diffusion coefficient D is made to vary with the concentration levels, resulting in a non-linear Fick's diffusion model given by,

$$\frac{dc}{dt} = \operatorname{div} \left(D(c) \operatorname{grad} c \right). \tag{4}$$

Here c denotes the relative concentration (rel. concentration) of the moisture in the material. The diffusion coefficient is a linear function of c, which is given by

$$D(c) = D_o + D_c c, \tag{5}$$

where D_o and D_c are the model parameters.

3 COUPLING MECHANICAL STIFFNESSES TO THE CONCENTRATION

The Lamé parameters for the dry and saturated material were determined with the help of relaxation experiments, where the dry material was modelled with the help of the Maxwell elements with longer relaxation times (j = 2, 3, 4) and the saturated material with the shorter relaxation times (j = 1, 2, 3) (Table 1). The Lamé parameters for any concentration c is calculated with the help of interpolating the

Maxwell element (j)	basic elasticity	j = 1	j = 2	j = 3	j = 4
Relaxation time [s]		0.5	5	50	500
Lamé Parameter (λ_{dry}) Dry [MPa]	1901.93	-	399.2	0.72	1.24
Lamé Parameter (λ_{wet}) Saturated [MPa]	270.83	359.03	446.61	172.41	-

Table 1: Lamé parameter λ for all Maxwell elements for dry and saturated material. Here λ_{dry} and λ_{wet} denote the vector of all the spring elements for dry and saturated material respectively. Since the Poisson's ratio for the material is unaffected from the moisture content, only λ is listed and μ can be calculated with the Poisson's ratio determined through experiments [21].

two extreme values of dry (λ_{dry}) and saturated (λ_{wet}) for all Maxwell elements (j = 1...4) including the basic elasticity. An interpolation function f(c) needs to be chosen so the interpolation,

$$\boldsymbol{\lambda}(c) = f(c) \, \boldsymbol{\lambda}_{dry} + (1 - f(c)) \, \boldsymbol{\lambda}_{wet} \tag{6}$$

can be done. The interpolation in [21] was done with a sigmoid function, fit with the help of just three points, which was not ideal. The change in the material properties happen because of a phase change from a glassy to rubbery phase. With more moisture, the glass transition temperature of the polymer reduces. This measurement was done by Sambale et al. [20] with the help of differential scanning calorimetry (DSC) and gives the variation of the glass transition temperature at different saturation levels. The measurement shows that the variation follows an exponential nature. The variation of the glass transition with saturation level was normalised in the range of zero to one, which is used to interpolate the Lamé parameters at a given concentration. An exponentially decreasing function,

$$f(c) = \exp(-5c) \tag{7}$$

was used to fit the curve between normalised glass transition temperature and the rel. concentration (Figure 2).

4 COUPLING CONCENTRATION TO THE APPLIED LOADING

The coupling of the diffusion model is generally done with the help of coupling the diffusion coefficient with the strain value. This is motivated from the fact that under tensile loading the free volume will increase and under compressive loading the free volume will decrease. The changing free volume will affect the speed of moisture uptake and hence the diffusion coefficient should change with the strain values. However in an experiment conducted at LKT, Technical University of Dortmund it was found out that the loading doesn't effect the diffusion speed.



Figure 2: (a) Glass transition temperature for different moisture uptake as weight percentages. (b) The glass transition temperature normalised and fit with decreasing exponential curve.



Figure 3: Loaded and unloaded material showing the same diffusion rate.

4.1 Diffusion under loading

An inhouse relaxation test jig was setup, where 12 different samples could be loaded simultaneously. Each of the 12 samples could be stored under sealed environmental chambers while loading and while maintaining the relative humidity (r.H.) level. To test the effect of loading on diffusion, two sets of samples were clamped under 100 % r.H., where one set of sample was stretched by a strain of 0.5 % and the other set was not loaded at all. The samples from both the sets were taken out regularly and at the same time (after 25 h, 50 h, 75 h, 100 h and 125 h) to weigh them. The sample taken out from the loaded set is not loaded again, rather another sample is taken out the next time to compare its weight with the unloaded sample. The plot of the mass uptake with time (Figure 3) shows that the moisture uptake speed for both sets of the samples were the same. This suggests that the diffusion process is not dependent on the mechanical loading. The assumption that the free volume changes and hence the diffusion speed changes may not hold, as the available integral free volume remains the same. However it is worth mentioning that the experiment was done for small deformation and at higher deformation, experiments at higher deformation were not conducted.

4.2 3-point-bending experiments

In order to validate the model under a moisture gradient 3-point bending (3PB) tests were conducted. The tests were conducted on 50 % saturated material. A specimen of the size 40 mm \times 4 mm \times 2 mm was used for the experiments, where the loading was applied on the face with the dimensions 40 mm \times 4 mm. In this way the concentration gradient on the specimen along the 2 mm edge was more dominant than the gradient across the 4 mm length (Figure 4). The beam was bent until 100 % of the



Figure 4: Load applied on the 40 mm \times 4 mm face, so the concentration gradient along the 2 mm edge is more dominant.



Figure 5: A comparison of the simulation and experimental results with one sided coupled model at 50 % saturation for (a) 0.1 mm/min and (b) 1 mm/min deflection rate.

thickness (2 mm) at two different rates of 0.1 mm/min and 1 mm/min. The tests were conducted in the dynamic mechanical analysis (DMA) Gabo Eplexor 500N machine (Netzsch, Deutschland). The results of the 3PB are compared to the simulation results of the one sided coupled model. The vertical force in the bending direction is used to compare the simulation and experimental results, as the stress value is highly dependent on the location of the neutral axis which can wander due to the inhomogeneity caused by moisture distribution.

Although the simulation can reproduce the results until 25 % of the total deflection, there is a notable deviation at higher deflection levels. The simulation results are less stiff than the experimental results as can be seen in Figure 5. This difference in the results are suspected to be a result of moisture redistribution due to the applied loading. Due to the pressure applied by the 3PB stamp on the beam, the moisture must be squeezed out and redistributed in the thickness. It can also be seen that for the slower deflection rate, the simulation shows more deviation than the faster rate. This shows that as the specimen gets more time under loading, the redistribution is more prominent. Thus, a coupling between the load applied and moisture transport does exist.



Figure 6: If the mass of water m_w in an element of volume *v* is subjected to loading, then the change in the volume to v_{new} is used to calculate the new concentration m_w/v_{new} .

The duration of the 3PB at the slowest deflection rate is just 1200 s. The process of diffusion is too slow to show an effect in this duration and the results in section 4.1 has shown that the diffusion process is independent of the loading. Hence, it can be postulated that the moisture redistribution occurs due to convection.

In order to simulate this behaviour, the change of the volume is used to model the transport. As the concentration is calculated by mass per volume, a transport or change of the volume results in a change of the concentration (Figure 6). This is done instead of introducing diffusion-convection equation for several reasons. Firstly, the numerical stability of a coupled model with a convection-diffusion equation is not reliable. Secondly, the flow of moisture due to convection cannot be restricted to a range of dry (c = 0) and fully saturated (c = 1) concentrations. Moreover, the velocity driving the convection has to be calculated with the help of another constitutive equation paralleling Darcy's law, thus making the model a 3-way coupled problem.

The volumetric strain (ε_v) is hence employed to calculate the new volume and the new rel. concentration c_{new} . The volumetric strain is given by,

$$\boldsymbol{\varepsilon}_{v} = \operatorname{tr} \boldsymbol{\varepsilon}. \tag{8}$$

and the new rel. concentration is given by,

$$c_{new} = \frac{c}{1 + \alpha \,\varepsilon_{\nu}},\tag{9}$$

where α is a model parameter. The correction to the rel. concentration is done after solving the set of equations for the one sided model in a seperate step. It doesn't effect the Jacobi matrix for the coupled model, but still introduces a change in the concentration due to loading.

5 RESULTS WITH COUPLING

The difference in the moisture distribution with and without coupling can be seen in Figure 7. With the coupled model, the volume of the upper part of the beam is compressed leading to higher rel. concentration, whereas in the one sided coupled model the distribution is symmetric along the beam thickness (Figure 7).

The force vs. deflection curve also shows a better correlation than the one without coupling (Figure 5). The reason for the better match is that a much larger portion of the width has lesser concentration than the one predicted by the one sided coupled model. Due to the decrease in the moisture content, the force value starts to increase after a certain amount of deflection. It can be seen that the force keeps on increasing with increasing deflection, since the concentration keeps on decreasing until it has no moisture along



Figure 7: The rel. concentration profile along the thickness of the testing specimen at the loading point. On the left with one sided coupled model and on the right with the fully coupled model.



Figure 8: A comparison of the simulation and experimental results with two sided coupled model at 50 % saturation for (a) 0.1 mm/min and (b) 1 mm/min deflection rate.

the width. This is also the case with the experimental results. However, in the case of the 1 mm/min simulation, this increase is more drastic than experiment, which can be altered with the parameter α . For a more accurate result, a convection term should be introduced, which as stated earlier leads to various numerical difficulties.

6 CONCLUSIONS

A material model that couples the mechanical deformation with the rel. concentration is presented here. A linear viscoelastic material model is used whose stiffness parameters are dependent on the rel. concentration of the moisture. The relationship between stiffness parameters and the rel. concentration is derived from the glass transition variation. The reverse coupling of the concentration on the loading is done by calculating the corrected volume and adjusting the rel. concentration according to the current volume. The volumetric part of the strain is used to calculate the current volume and hence the concentration.

3-point-bending tests were used to validate the model. A comparison of the one sided coupled model with the experimental result shows that the simulation can match the experiment up to a certain value but with an increase in the deflection the experiment gives higher force values or a stiffer material. This can be attributed to the loss of moisture with the loading. With a two way coupled model the simulation can represent the experimental values more accurately. There is scope for improvement of the material model by introducing a pressure term that comes from the stress applied on the material and which induces a transport of the moisture already absorbed by the material.

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