COUPLED ADJOINT FLUID-STRUCTURE INTERACTION TECHNIQUE FOR FLEXIBLE WING SHAPE OPTIMIZATION

K. TSIAKAS¹, G. ROGÉ², X. TROMPOUKIS¹, V. ASOUTI¹, G. FOUGERON³, S. JULISSON², L. MARTIN², S. KLEINVELD² AND K. GIANNAKOGLOU¹

¹ Parallel CFD & Optimization Unit (PCOpt), School Of Mechanical Engineering, National Technical University of Athens (NTUA), 9, Iroon Polytechniou, 15780, Zografou, Greece {tsiakost, xeftro}@gmail.com, {vasouti, kgianna}@mail.ntua.gr

² Dassault Aviation 78, quai Marcel Dassault, 92552, Saint-Cloud, France {gilbert.roge, sarah.julisson, ludovic.martin, steven.kleinveld}@dassault-aviation.com

> ³ ESI Group, 3 Rue Saarinen, 94150 Rungis, France gabriel.fougeron@esi-group.com

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Abstract. This paper presents the Multi-Disciplinary Optimization (MDO) of a business jet including static aeroelastic effects. Two different CFD tools (AETHER by Dassault Aviation and PUMA by NTUA) are coupled with a CSM model (VPS software by ESI). Single discipline as well as coupled multi-disciplinary sensitivity derivatives (SDs) are computed by means of adjoint methods. Both a purely discrete and a hybrid continuous(fluid)/discrete(structure) adjoint formulations are presented. The computed SDs are verified against finite differences.

1 INTRODUCTION

Modern aircraft design is an increasingly complex process due to stricter regulations, environmental constraints and a continuous need for lower operational costs. Design includes aerodynamics, structural mechanics, control, aeroacoustics and heat transfer, to name just the most important disciplines. There is a large literature on simulation tools for these disciplines, the use of which has led to efficient and increasingly versatile aircraft. A major challenge for future design workflows, especially those relying on numerical optimization techniques, is the exploitation of trade-offs between the involved disciplines which is expected to reduce the development effort, resulting in superior new designs.

Among the variety of optimization techniques used by the aeronautics industry, a significant amount of interest has surfaced in gradient-based techniques relying upon the adjoint method [1, 2, 3], since the latter computes gradients for aircraft performance metrics at a cost independent of the number of design parameters. However, adjoint formulations that account for the interaction between disciplines, are not yet mature enough and this is where recent research focuses on [4, 5]. The development of adjoint methods in an MDO context gives the opportunity to obtain feasible designs at a fraction of the development cost compared to traditional methods, iteratively solving single-discipline optimization problems. The latter are more costly yielding inferior results, since the multi-disciplinary feasibility must

a posteriori be enforced for each candidate design.

In this work, adjoint methods for tackling the problem of aircraft wing design while also considering the wing's structure flexibility are developed. These make use of coupled high-fidelity computational tools (for the Reynolds-Averaged Navier-Stokes, RANS, equations for the fluid and a Finite Element, FE, solver for the wing structure) to support flexible wing optimization. The developed methods are used for the shape optimization of a Generic Business Jet (GBJ) flexible wing provided by Dassault Aviation.

The rest of this paper is as follows. Section 2 describes the two different CFD solvers, namely AETHER and PUMA, used in this paper. Comparisons of the GBJ aerodynamic performance predictions by them are made. Section 3 describes the adjoint formulations for both CFD tools. AETHER employs discrete adjoint based on Automatic Differentiation (AD), while PUMA continuous adjoint. Section 4 describes the approach used for parameterizing the shape of the GBJ. A surrogate parameterization model provides an opaque layer of abstraction between the actual CAD model of the geometry and the optimization process. Section 5 describes the Computational Structural Mechanics (CSM) model used, by the VPS software, as well as the Fluid-Structure Interaction (FSI) technique for both the primal and adjoint workflows. In Section 6, the sensitivity derivatives predicted by the adjoint methods, with and without accounting for wing structure flexibility,0 are compared with finite differences (FDs). Section 7 presents the outcome of the optimization runs.

2 CFD TOOLS

Two CFD tools are used. The first one is code AETHER, namely the in-house software of Dassault Aviation. AETHER is based on the Streamline Upwind Petrov-Galerkin (SUPG) Finite Element approach for solving the RANS equations on simplicial meshes [6, 7]. A fully implicit iterative timemarching procedure based on the GMRES algorithm leads to a robust and efficient solution scheme. AETHER runs on CPU clusters using domain decomposition and the MPI protocol.

The second CFD tool is software PUMA [8] developed by the PCOpt/NTUA. PUMA is a GPUenabled (U)RANS solver, employing the vertex-centered approach of the finite volume method on unstructured meshes consisting of tetrahedra, pyramids, prisms and hexahedra. Inviscid fluxes are discretized using either Roe's upwind approximate Riemann solver or a central scheme with a blend of second- and fourth-order dissipation terms for second-order spatial accuracy. It employs a multi-stage Runge-Kutta scheme with residual smoothing, for which the linear system is solved by a preconditioned Krylov-based solver. PUMA runs on clusters of GPUs on different computational nodes, by decomposing the flow domain into overlapping subdomains. GPUs on the same node exchange data through the shared on-node memory using the capabilities of the CUDA API, while the MPI protocol is used otherwise.

A comparison of the c_L vs. c_D curves for a cruise speed of M = 0.82 and a flight altitude of 41 000 ft, as computed by the two CFD tools, is presented in Fig. 2. The two curves match each other and the comparison is very satisfactory given that the two RANS solvers rely upon different methods (FE vs. FV) and meshes. AETHER used a tetrahedral mesh created by in-house Dassault Aviation tools with approximately 16 Mi nodes while PUMA used a hybrid mesh of 9 Mi nodes, Fig. 1. In either code, turbulence is modeled by means of the Spalart-Allmaras model [9].



Figure 1: Cut across the volume mesh around the GBJ used by PUMA (left). Hexahedra and prisms are used for creating high aspect ratio cells to capture the boundary layer physics. Tetrahedra fill the rest of the domain, while pyramids are used to transition between the different element types. Close-up view of the mesh in the nacelle-pylon-fuselage area (right).



Figure 2: Comparison of the c_L vs. c_D curves predicted by AETHER and PUMA.

3 ADJOINT FORMULATIONS

The two CFD tools (PUMA and AETHER) employ different adjoint approaches. In the discrete adjoint approach, the expression of the objective or constraint function J is augmented by the discrete residual of the flow (index F) equations \mathbf{R}_n^F , multiplied by the vector of adjoint variables Ψ_n^F , with $n=1,\ldots,M^F$, with M^F the number of the flow equations. Boldface symbols denote vectors with size equal to that of the computational mesh. By differentiating the expression of J_{aug} w.r.t. the design parameters b_i , $i=1,\ldots,N^D$ the following expression arises

$$\frac{\delta J_{\text{aug}}}{\delta b_i} = \frac{\delta J}{\delta b_i} + [\boldsymbol{\Psi}_n^F]^T \frac{\delta \boldsymbol{R}_n^F}{\delta b_i} \tag{1}$$

Throughout this paper, repeated indices imply summation, unless otherwise stated. Since both J and \mathbf{R}_n^F are functions of the flow variables \mathbf{W}_m^F , $m = 1, ..., M^F$ and nodal coordinates \mathbf{x}_k^F , k = 1, ..., 3, their

variations w.r.t. b_i are

$$\frac{\delta J}{\delta b_i} = \frac{\partial J}{\partial \boldsymbol{W}_m^F} \frac{\delta \boldsymbol{W}_m^F}{\delta b_i} + \frac{\partial J}{\partial \boldsymbol{x}_k^F} \frac{\delta \boldsymbol{x}_k^F}{\delta b_i}, \qquad \frac{\delta \boldsymbol{R}_n^F}{\delta b_i} = \frac{\partial \boldsymbol{R}_n^F}{\partial \boldsymbol{W}_m^F} \frac{\delta \boldsymbol{W}_m^F}{\delta b_i} + \frac{\partial \boldsymbol{R}_n^F}{\partial \boldsymbol{x}_k^F} \frac{\delta \boldsymbol{x}_k^F}{\delta b_i}$$
(2)

By combining Eq. 1 with Eq. 2 and re-arranging terms, one obtains

$$\frac{\delta J_{\text{aug}}}{\delta b_i} = \left(\frac{\partial J}{\partial \boldsymbol{W}_m^F} + [\boldsymbol{\Psi}_n^F]^T \frac{\partial \boldsymbol{R}_n^F}{\partial \boldsymbol{W}_m^F}\right) \frac{\delta \boldsymbol{W}_m^F}{\delta b_i} + \left(\frac{\partial J}{\partial \boldsymbol{x}_k^F} + [\boldsymbol{\Psi}_n^F]^T \frac{\partial \boldsymbol{R}_n^F}{\partial \boldsymbol{x}_k^F}\right) \frac{\delta \boldsymbol{x}_k^F}{\delta b_i}$$
(3)

Variations in the flow variables are eliminated by satisfying the discrete adjoint equations (i.e. $\left[\frac{\partial R_n^F}{\partial W_m^F}\right]^T \Psi_n^F = -\left[\frac{\partial J}{\partial W_m^F}\right]^T$), while the second term on the r.h.s. of Eq. 3 provides the sensitivity derivatives' expression.

The discrete adjoint approach, used by AETHER [10] is based on the AD tool TAPENADE [11]. The Spalart-Allmaras turbulence model is also included in the AD process, hence, \mathbf{R}^F also accommodates the residual of the discretized Spalart-Allmaras model equation.

In order to compute SDs using continuous adjoint, J is augmented by the integrals of the flow equations R^F multiplied by the adjoint variables Ψ^F . No boldface symbols are used since the adjoint development is agnostic to any domain discretization. Development starts from

$$\frac{\delta J_{\text{aug}}}{\delta b_i} = \frac{\delta J}{\delta b_i} + \underbrace{\int_{\Omega^F} \Psi_n^F \frac{\delta R_n^F}{\delta b_i} d\Omega}_{\mathcal{T}_1} + \underbrace{\int_{\Omega^F} \Psi_\ell^G \frac{\partial^2}{\partial x_k^2} \left(\frac{\delta x_\ell^F}{\delta b_i}\right) d\Omega}_{\mathcal{T}_2}$$
(4)

where $n = 1, ..., M^F$, $k, \ell = 1, ..., 3$ and Ω^F is the CFD domain. A Laplacian mesh deformation model is assumed and this is also included in J_{aug} through \mathcal{T}_2 ; this introduces three extra adjoint fields Ψ_m^G . By applying the divergence theorem on \mathcal{T}_1 and \mathcal{T}_2 and further developing $\frac{\delta J}{\delta b_i}$ one obtains

$$\frac{\delta J}{\delta b_{i}} = \int_{S_{Obj}} \mathcal{B}_{m}^{J \to F} \frac{\delta W_{m}^{F}}{\delta b_{i}} dS + \int_{S_{Obj}} \mathcal{B}_{k}^{J \to G} \frac{\delta x_{k}^{F}}{\delta b_{i}} dS$$

$$\mathcal{T}_{1} = \int_{\Omega^{F}} \mathcal{C}_{m}^{F \to F} \frac{\delta W_{m}^{F}}{\delta b_{i}} d\Omega + \int_{\Omega^{F}} \mathcal{C}_{k}^{F \to G} \frac{\delta x_{k}^{F}}{\delta b_{i}} d\Omega + \int_{S} \mathcal{B}_{m}^{F \to F} \frac{\delta W_{m}^{F}}{\delta b_{i}} dS + \int_{S} \mathcal{B}_{k}^{F \to G} \frac{\delta x_{k}^{F}}{\delta b_{i}} dS \qquad (5)$$

$$\mathcal{T}_{2} = \int_{\Omega^{F}} \mathcal{C}_{k}^{G \to G} \frac{\delta x_{k}^{F}}{\delta b_{i}} d\Omega + \int_{S} \mathcal{B}_{k}^{G \to G} \frac{\delta x_{k}^{F}}{\delta b_{i}} dS$$

where S stands for the boundary of Ω , S_{Obj} the surface on which J is defined, $m = 1, ..., M^F$ and k = 1, ..., 3. Combining Eqs. 5 and satisfying (a) the adjoint flow equations $C_m^{F \to F} = 0$ with appropriate boundary conditions, i.e. $\mathcal{B}_m^{J \to F} + \mathcal{B}_m^{F \to F} = 0$ and (b) the adjoint grid displacement equations $C_k^{F \to G} + C_k^{G \to G} = 0$, the derivatives of J w.r.t. b_i are computed as

$$\frac{\delta J}{\delta b_i} = \int\limits_{S} \left(\mathcal{B}_k^{J \to G} + \mathcal{B}_k^{F \to G} + \mathcal{B}_k^{G \to G} \right) \frac{\delta x_k^F}{\delta b_i} \mathrm{d}S \tag{6}$$

In the adjoint PUMA code, the Spalart-Allmaras turbulence model is fully differentiated [12, 13]. Possible numerical stability issues of the adjoint solver due to the presence of unstable modes in the primal solution, are circumvented using the Recursive Projection Method (RPM) [14].

4 GBJ SHAPE PARAMETERIZATION

The shape of the airframe is parameterized using the geometry generation tool GANIMEDE (Geometry ANd Inherent MEsh DEformation, [15, 16]), by Dassault Aviation. GANIMEDE handles both local and global design variables. By local design variables modifications at the control point level are meant, such as coordinates', tangent and curvature values. Global design variables redefine several control points enabling to modify characteristics such as thickness, twist and camber of wing sections. As the geometry changes during the optimization, GANIMEDE generates new surface meshes by involving the CAD modeler and creating connectivity by projecting the initial surface mesh onto the new geometry. Quite often, the need to take geometric constraints into account emerges occasionally. For manufacturing reasons, one might want to impose specific relations between some geometric entities. From the operational point of view, for many aircraft, including very large commercial transport concepts, a span constraint has to be imposed in the design process. Another case of interest for imposing constraints exists when retrofitting an aircraft. All geometric constraints are directly handled by the GANIMEDE modeler. This results to a cost reduction compared to the alternative of delegating this task to the optimizer by imposing additional constraints, since the gradient of these geometric quantities is no longer needed.

Since GANIMEDE is not available for use outside Dassault Aviation, NTUA has developed a surrogate shape parameterization model, working directly on the aircraft's surface mesh, based on meshes corresponding to shape perturbations for each design parameter, provided using GANIMEDE. By doing so, the dependence on GANIMEDE is restricted only to the pre-processing phase, while shape modifications during the optimization are undertaken by the surrogate model.

Let z^D be the coordinates of the surface mesh nodes for the datum shape and $b_i \in [b_i^-, b_i^+]$, $i = 1, ..., N^D$ the design parameters with initial (datum) value of b_i^D . Let z_i^- , $i = 1, ..., N^D$ be the surface mesh nodes coordinates corresponding to the design parameter vector with all parameters equal to their datum values, excluding though the *i*th one which is set to b_i^- , also, let z_i^+ be the coordinates for the positive perturbation of the *i*th design parameter. The z_i values for any value of the *i*th design parameter are given by

$$\mathbf{z}_{ik} = \mathbf{z}_k^D + \mathbf{a}_{ik}b_i^2 + \mathbf{c}_{ik}b_i$$
, with $i = 1, \dots, N^D$ and $k = 1, \dots, 3$ (7)

(no summation for repeated indices). The coefficients a and c are computed by solving the following system of equations for each surface mesh node, design parameter and Cartesian coordinate

$$\boldsymbol{a}_{ik} \left(\boldsymbol{b}_{i}^{-}\right)^{2} + \boldsymbol{c}_{ik} \boldsymbol{b}_{i}^{-} = \boldsymbol{z}_{ik}^{-} - \boldsymbol{z}_{k}^{D}$$

$$\boldsymbol{a}_{ik} \left(\boldsymbol{b}_{i}^{+}\right)^{2} + \boldsymbol{c}_{ik} \boldsymbol{b}_{i}^{+} = \boldsymbol{z}_{ik}^{+} - \boldsymbol{z}_{k}^{D}$$
(8)

The surface mesh nodal coordinates x for a simultaneous modification of several design parameters results from superimposing the various z values for all $i = 1, ..., N^D$, i.e.

$$\mathbf{x}_{k} = \mathbf{z}_{k}^{D} + \sum_{i=1}^{N^{D}} (\mathbf{z}_{ik} - \mathbf{z}_{k}^{D})$$
, for $k = 1, ..., 3$ (9)

In what follows, 4 design parameters $(b_1 \text{ to } b_4)$ are used to control the wing twist spanwise distribution, another 4 $(b_5 \text{ to } b_8)$ to control the wing trailing edge camber distribution along its span and the last one (b_9) to control the horizontal tail plane (HTP) rotation angle. Since this model is analytically differentiated, it can readily be incorporated into the adjoint method.



Figure 3: Effect of the design parameters on the airframe shape. The datum contour of the wing is shown with the blue line. The yellow line corresponds to the altered shape outline. Alteration of the 1st (bottom-right), 8th (top-right) and 9th (bottom-left) design parameters.

5 FLUID-STRUCTURE INTERACTION (FSI)

The purpose is to introduce wing flexibility effects into the analysis and optimization of an aircraft. The CFD tools and their adjoint counterparts represent only one discipline involved in the MDO problem. The other discipline concerns the structural response of the wing under aerodynamic and gravitational loading. This is handled by a FE linear elasticity solver, part of the VPS (Virtual Performance Solution) software of ESI, which is the CSM component of the MDO workflow presented.

Given the positions \mathbf{x}_k^S of the N^S FE model nodes on the flight shape, and denoting by \mathbf{y}_k^S the positions of the same nodes in the unloaded state (jig shape), the nodal structural displacements are given as $\mathbf{U}_k^S = \mathbf{x}_k^S - \mathbf{y}_k^S$, k = 1, ..., 3. The discrete system of the linear elasticity equations along the *k* direction is expressed as

$$\mathbb{K}_{k\ell} \boldsymbol{U}_{\ell}^{S} = \boldsymbol{f}_{k}^{S} + \boldsymbol{f}_{k}^{G}, \quad \text{with} \quad k, \ell = 1, \dots, 3$$

$$(10)$$

where \mathbb{K} stands for the stiffness matrix, f_k^S for the aerodynamic load along the Cartesian direction k applied on the structure nodes and f_k^G for the gravitational load due to the structure weight (including also the lumped fuel load). Starting from the wing's flight shape, a FE model, consisting of shell elements (for the wing's skin and webs of spars and ribs) and beam elements (for the caps of spars and ribs), is set up. Then, the jig shape of the wing is computed iteratively by a quasi-Newton method, assuming \mathbb{K} to be a function of y^S .

The CSM mesh differs from the surface CFD mesh. The tool for transferring information along the non-matching interfaces has been developed by ESI, based on Radial Basis Functions (RBFs) [17]. The structural displacements of the CFD surface mesh nodes are given as $U_k^F = \mathbb{H}U_k^S$, k = 1,...,3, where matrix \mathbb{H} contains the weights with which each of the N^S nodes of the CSM model contributes to the displacement of the N^F CFD surface mesh nodes. Aerodynamic load transfer along the FSI is done using the transpose of \mathbb{H} , and this ensures conservation of force, moment and virtual work along the FSI [18].

The aeroelastic analysis of the aircraft is performed using a fixed-point iteration scheme with dynamic relaxation [19], sketched in Fig. 4. In order to compute sensitivity derivatives of the aerodynamic performance metrics, while accounting for wing flexibility, the adjoint technique of Section 3 has been extended. Eq. 1 now includes an additional term containing variations in the discrete CSM residual $(\frac{\delta R^{S}}{\delta h})$ multiplied by the CSM adjoint variables Ψ^S , i.e.



Figure 4: The aeroelastic fixed-point iteration involving the CFD and CSM tools.

$$\frac{\delta J_{\text{aug}}}{\delta b_i} = \frac{\delta J}{\delta b_i} + [\boldsymbol{\Psi}_n^F]^T \frac{\delta \boldsymbol{R}_n^F}{\delta b_i} + [\boldsymbol{\Psi}_k^S]^T \frac{\delta \boldsymbol{R}_k^S}{\delta b_i}, \text{ with } n = 1, \dots, M^F \text{ and } k = 1, \dots, 3.$$
(11)

In this work, the wing structure is assumed to remain unaffected by changes in the design parameters. Moreover, the wing's datum flight shape (with nodal coordinates \mathbf{x}^{F_0}) is parameterized and will be referred to onward as the "reference" shape. Let \mathbf{y}_k^F denote the coordinates of the wing skin nodes in the unloaded (jig shape) state. Then, the actual flight shape of the wing (when static aeroelastic equilibrium is obtained) is given as $\mathbf{x}_k^F = \mathbf{y}_k^F + \mathbf{U}_k^F$, k = 1, ..., 3. Due to the assumption made, $\mathbf{d}_k^F = \mathbf{x}_k^{F_0} - \mathbf{y}_k^F$, k = 1, ..., 3 remain constant and the variation of the flight shape coordinates reads

$$\frac{\delta \mathbf{x}_{k}^{F}}{\delta b_{i}} = \frac{\delta \mathbf{x}_{k}^{F_{0}}}{\delta b_{i}} + \frac{\delta \mathbf{U}_{k}^{F}}{\delta b_{i}}, \text{ with } k = 1, \dots, 3$$
(12)

The objective or constraint function J, the aerodynamic loads on the CFD surface mesh nodes f^F and the residuals of the flow equations R^F are now functions of the flow variables W^F , the reference shape nodal coordinates x^{F_0} and the displacements induced by the wing flexibility U^F . Moreover, the residuals of the CSM equations R^S depend on the structural displacement of the CSM nodes U^S and the aerodynamic loads transferred to the CSM model f^S . Eq. 11 then, becomes

$$\frac{\delta J_{\text{aug}}}{\delta b_{i}} = \left(\frac{\partial J}{\partial \boldsymbol{W}_{m}^{F}} + [\boldsymbol{\Psi}_{n}^{F}]^{T} \frac{\partial \boldsymbol{R}_{n}^{F}}{\partial \boldsymbol{W}_{m}^{F}}\right) \frac{\delta \boldsymbol{W}_{m}^{F}}{\delta b_{i}} + [\boldsymbol{\Psi}_{k}^{S}]^{T} \frac{\partial \boldsymbol{R}_{k}^{S}}{\partial \boldsymbol{U}_{\ell}^{S}} \frac{\delta \boldsymbol{U}_{\ell}^{S}}{\delta b_{i}} + \underbrace{\left(\frac{\partial J}{\partial \boldsymbol{U}_{\ell}^{F}} + [\boldsymbol{\Psi}_{n}^{F}]^{T} \frac{\partial \boldsymbol{R}_{n}^{F}}{\partial \boldsymbol{U}_{\ell}^{F}}\right) \frac{\delta \boldsymbol{U}_{\ell}^{F}}{\delta b_{i}}}_{\mathcal{T}_{3}} + \underbrace{\left(\boldsymbol{\Psi}_{k}^{S}\right)^{T} \frac{\partial \boldsymbol{R}_{k}^{S}}{\partial \boldsymbol{f}_{\ell}^{S}} \frac{\delta \boldsymbol{f}_{\ell}^{S}}{\delta b_{i}}}_{\mathcal{T}_{4}} + \left(\frac{\partial J}{\partial \boldsymbol{x}_{\ell}^{F_{0}}} + [\boldsymbol{\Psi}^{F}]^{T} \frac{\partial \boldsymbol{R}_{n}^{F}}{\partial \boldsymbol{x}_{\ell}^{F_{0}}}\right) \frac{\delta \boldsymbol{x}_{\ell}^{F_{0}}}{\delta b_{i}}}{(13)}$$

with $m, n = 1, ..., M^F$ and $k, \ell = 1, ..., 3$. The multiplier of the variation of U^F is the so-called adjoint structural displacements vector $(U^{F,adj})$ on the CFD surface mesh nodes, while that of the variation of f^S represents the adjoint aerodynamic loads $(f^{S,adj})$ on the CSM model nodes. Taking this into account, as well as the equations connecting U^F to U^S and f^S to f^F , terms \mathcal{T}_3 and \mathcal{T}_4 in Eq. 13 become

$$\mathcal{T}_{3} = \underbrace{[\boldsymbol{U}_{m}^{F,\mathrm{adj}}]^{T}}_{[\boldsymbol{U}_{\ell}^{S,\mathrm{adj}}]^{T}} \frac{\partial \boldsymbol{U}_{m}^{F}}{\partial \boldsymbol{U}_{\ell}^{S}} \frac{\delta \boldsymbol{U}_{\ell}^{S}}{\delta b_{i}}, \qquad \mathcal{T}_{4} = \underbrace{[\boldsymbol{f}_{k}^{S,\mathrm{adj}}]^{T}}_{[\boldsymbol{f}_{q}^{F,\mathrm{adj}}]^{T}} \frac{\partial \boldsymbol{f}_{k}^{S}}{\partial \boldsymbol{f}_{q}^{F}} \frac{\delta \boldsymbol{f}_{q}^{F}}{\delta b_{i}}$$
(14)

The final expression of the discrete adjoint equations for the coupled CFD+CSM system of equations is

$$[\boldsymbol{\Psi}_{n}^{F}]^{T}\frac{\partial\boldsymbol{R}_{n}^{F}}{\partial\boldsymbol{W}_{m}^{F}} + [\boldsymbol{f}_{k}^{F,\mathrm{adj}}]^{T}\frac{\partial\boldsymbol{f}_{k}^{F}}{\partial\boldsymbol{W}_{m}^{F}} = -\frac{\partial J}{\partial\boldsymbol{W}_{m}^{F}}, \qquad [\boldsymbol{\Psi}_{k}^{S}]^{T}\frac{\partial\boldsymbol{R}_{k}^{S}}{\partial\boldsymbol{U}_{\ell}^{S}} = -[\boldsymbol{U}^{S,\mathrm{adj}}]^{T} - [\boldsymbol{f}_{q}^{F,\mathrm{adj}}]^{T}\frac{\partial\boldsymbol{f}_{q}^{F}}{\partial\boldsymbol{U}_{k}^{F}}\frac{\partial\boldsymbol{U}_{k}^{F}}{\partial\boldsymbol{U}_{\ell}^{S}} \quad (15)$$

with $m, n = 1, ..., M^F$ and $k, \ell, q = 1, ..., 3$. Eqs. 15 are solved using the same iterative scheme used for the primal aeroelastic solution. Once Ψ^F and Ψ^S have been computed, the sensitivity derivatives of J w.r.t. b_i are obtained as

$$\frac{\delta J}{\delta b_i} = \left(\frac{\partial J}{\partial \boldsymbol{x}_k^{F_0}} + [\boldsymbol{\Psi}_m^F]^T \frac{\partial \boldsymbol{R}_m^F}{\partial \boldsymbol{x}_k^{F_0}} + [\boldsymbol{f}_\ell^{F, \mathrm{adj}}]^T \frac{\partial \boldsymbol{f}_\ell^F}{\partial \boldsymbol{x}_k^{F_0}}\right) \frac{\delta \boldsymbol{x}_k^{F_0}}{\delta b_i}$$
(16)

with $m = 1, ..., M^F$ and $k, \ell = 1, ..., 3$, where the variations in $\mathbf{x}_k^{F_0}$ are computed by analytically differentiating Eq. 9.

In the continuous adjoint of PUMA, the aeroelastic adjoint development is largely similar. Eq. 4 is further extended by a discrete part containing the structural adjoint variables and variation in the residuals of the discrete CSM equations ("hybrid" continuous-discrete FSI adjoint), i.e.

$$\frac{\delta J_{\text{aug}}}{\delta b_i} = \frac{\delta J}{\delta b_i} + \underbrace{\int_{\Omega^F} \Psi_n^F \frac{\delta R_n^F}{\delta b_i} d\Omega}_{\mathcal{T}_1} + \underbrace{\int_{\Omega^F} \Psi_m^G \frac{\partial^2}{\partial x^{F_k^2}} \left(\frac{\delta x_m^F}{\delta b_i}\right) d\Omega}_{\mathcal{T}_2} + [\Psi_m^S]^T \frac{\delta R_m^S}{\delta b_i}$$
(17)

The last term is processed as in the previous case. The first three terms of the r.h.s. of Eq. 17 are processed by taking into account Eq. 12. Thus, the corresponding equations read

$$\frac{\delta J}{\delta b_{i}} = \int_{S_{Obj}} \mathcal{B}_{k}^{J \to F} \frac{\delta W_{m}^{F}}{\delta b_{i}} dS + \int_{S_{Obj}} \mathcal{B}_{k}^{J \to G} \frac{\delta x_{k}^{F_{0}}}{\delta b_{i}} dS + \int_{S_{Obj}} \mathcal{B}_{k}^{J \to G} \frac{\delta U_{k}^{F}}{\delta b_{i}} dS$$

$$\mathcal{T}_{1} = \int_{\Omega^{F}} \mathcal{C}_{m}^{F \to F} \frac{\delta W_{m}^{F}}{\delta b_{i}} d\Omega + \int_{\Omega^{F}} \mathcal{C}_{k}^{F \to G} \frac{\delta x_{k}^{F}}{\delta b_{i}} d\Omega + \int_{S} \mathcal{B}_{m}^{F \to F} \frac{\delta W_{m}^{F}}{\delta b_{i}} dS + \int_{S} \mathcal{B}_{k}^{F \to G} \frac{\delta x_{k}^{F_{0}}}{\delta b_{i}} dS + \int_{S} \mathcal{B}_{k}^{F \to G} \frac{\delta x_{k}^{F}}{\delta b_{i}} dS + \int_{S} \mathcal{B}_{k}^{F \to G} \frac{\delta x_{k}^{F}}{\delta b_{i}} dS + \int_{S} \mathcal{B}_{k}^{F \to G} \frac{\delta x_{k}^{F}}{\delta b_{i}} dS + \int_{S} \mathcal{B}_{k}^{F \to G} \frac{\delta x_{k}^{F}}{\delta b_{i}} dS + \int_{S} \mathcal{B}_{k}^{F \to G} \frac{\delta x_{k}^{F}}{\delta b_{i}} dS + \int_{S} \mathcal{B}_{k}^{F \to G} \frac{\delta x_{k}^{F}}{\delta b_{i}} dS$$

$$\mathcal{T}_{2} = \int_{\Omega^{F}} \mathcal{C}^{G \to G} \frac{\delta x_{k}^{F}}{\delta b_{i}} d\Omega + \int_{S} \mathcal{B}_{k}^{G \to G} \frac{\delta x_{k}^{F_{0}}}{\delta b_{i}} dS + \int_{S} \mathcal{B}_{k}^{G \to G} \frac{\delta U_{k}^{F}}{\delta b_{i}} dS$$
(18)

Integrals containing variations in U^F are discretized and then combined with those containing variations of U^F arising from the last term of Eq. 17, effectively leading to the discrete expression of the adjoint structural displacements on the CFD surface mesh. The adjoint boundary conditions on the CFD domain are also different, since now the third integral in the second of equations Eq. 18 has to be discretized and, then, combined with terms arising from the variations of f^S hidden in the variations of R^S . Thus, the adjoint field equations are exactly the same as in the single-discipline case and the only difference (from the adjoint to the CFD point of view) is the new discrete expression of adjoint boundary conditions. Compared to the single-discipline case, the final expression of SDs is augmented only by a term same as the last term of Eq. 16.

6 VERIFICATION OF SENSITIVITY DERIVATIVES

The SDs computed by the adjoint of AETHER and PUMA, with and without aeroelastic effects, are compared to central FDs. Moreover, in order to quantify the accuracy loss due to the commonly used "frozen" turbulence assumption, the adjoint sensitivities are also computed by neglecting variations in turbulent quantities w.r.t. the design parameters. It is shown that, when including the turbulence model in the adjoint, the accuracy of SDs is significantly improved for both CFD tools and, hence, this approach is followed in all optimization runs, see Fig. 5.



Figure 5: SDs of c_L w.r.t. (a) design parameters controlling the wing twist distribution (and the HTP rotation) and (b) design parameters controlling the wing's trailing edge camber distribution along its span. The SDs computed by the adjoint of PUMA (continuous) and AETHER (discrete) are compared with FDs.

7 OPTIMIZATION RESULTS

Two optimization runs have been performed for optimizing the GBJ geometry, by considering or neglecting wing structure flexibility. The shape parameterization model is the one described in Section 4. Flow conditions correspond to $M_{inf}=0.82$, angle of attack of 2.5° and a flight altitude of 41 000 ft. In both cases, the SLSQP algorithm [20] is used to drive the search of the design space and PUMA is used as the primal and adjoint CFD tool. Both optimization runs target in minimizing the drag coefficient (c_D), with the lift coefficient (c_L) greater than or equal to that of the datum geometry (with a 0.01 threshold) and the absolute value of the pitching moment coefficient ($|c_M|$) less than 0.001, to get a trimmed configuration. The convergence history of the rigid wing optimization is presented in Fig. 6a. In the second optimization run, the structural response of the wing structure was taken into account in the evaluation of both the aerodynamic performance metrics and the corresponding sensitivity derivatives, by means of the coupled adjoint workflow of Section 5. The hybrid continuous(flow)-discrete(structure) approach employed in



the PUMA + VPS workflow is used. The optimization convergence history is presented and results are commented in Fig. 6b.

Figure 6: (a) Convergence of the rigid wing optimization. A solution satisfying both aerodynamic constraints is obtained after 12 cycles with ~ 5% reduction in c_D , compared to the datum shape. Reevaluating this design with an aeroelastic tool, the c_D reduction was 4.7% and the pitching moment constraint was violated. (b) Convergence of the flexible wing optimization. A solution satisfying both constraints was obtained after 15 optimization cycles with ~ 3% reduction in c_D , compared to the datum shape. Since, the wing structural response was taken into account, this is a feasible design w.r.t. aeroelastic criteria.

The design parameters corresponding to the rigid optimum are compared to those of the flexible optimum in Fig. 7. It is seen that, when flexibility is taken into account, the twist parameters are not varied significantly, since large variations lead to constraint violation. In Fig. 8, the pressure coefficient fields on the suction side of the datum, single-discipline optimum and multi-disciplinary optimum are compared. For the former, drag reduction is mostly achieved by reducing the shock intensity. This not so evident in the latter. Moreover, it can be seen that the loading of the wing at the winglet area is significantly increased for the flexible optimum.

8 CONCLUSIONS

The development of the adjoint method for computing aeroelastic sensitivities of a business jet flexible wing was presented. Two CFD codes (AETHER and its discrete adjoint and PUMA and its continuous adjoint) were coupled with a CSM FE model for the wing structure (VPS). Optimization runs show that neglecting aeroelastic effects during the optimization leads to potentially infeasible designs. Thus, the importance of the MDO approach for this class of optimization problems is convincingly demonstrated. The cost of computing coupled SDs is roughly 3x the cost of performing a single-discipline aerodynamic analysis and sensitivity derivative computation. Differences caused by the "frozen" turbulence assumption in an adjoint code are also demonstrated. On-going work focuses on the extension of the MDO



Figure 7: Comparison of the optimal set of (a) twist (parameters 1,2,3,4 and 9) and (b) trailing edge camber (parameters 5,6,7,8) design parameters located from the rigid and flexible optimizations.



Figure 8: Comparison of the pressure coefficient iso-areas along the suction side surface of the wing of the (a) datum, (b) rigid and (c) flexible optimal airframe.

framework to include variations in the structural model properties into the optimization process.

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